## Preserving invariants and volume for split systems

Philippe Chartier ${ }^{1} \quad$ Ander Murua ${ }^{2}$
${ }^{1}$ IPSO
INRIA-Rennes and ENS Cachan, Antenne de Bretagne
${ }^{2}$ Department of Computer Science
University of the Basque Country
Clermont 2008

## Outline

## (1) Problems and motivations

- General invariants encountered in physics
- Improved qualitative behavior of geometric integrators

2) Setting of the problem

- Invariant and volume preservation for split systems
- The Hopf algebra of coloured trees
- B -series and S -series for split vector fields
(3) Conditions for invariants-preservation
- Numerical methods preserving invariants
- The case of quadratic and cubic invariants
- B-series methods preserving all cubic invariants

4 Conditions for volume-preservation

- Volume-preserving B-series
- Connection with the preservation of cubic invariants
- volume preserving methods for split systems with a special stru
(5) From conditions for vector fields to conditions for integrators


## Outline

## (1) Problems and motivations

- General invariants encountered in physics
- Improved qualitative behavior of geometric integrators

2

## Setting of the problem

- Invariant and volume preservation for split systems
- The Hopf algebra of coloured trees
- B-series and S-series for split vector fieldsConditions for invariants-preservation
- Numerical methods preserving invariants
- The case of quadratic and cubic invariants
- B-series methods preserving all cubic invariants

Conditions for volume-preservation

- Volume-preserving B-series
- Connection with the preservation of cubic invariants
- volume preserving methods for split systems with a special stru From conditions for vector fields to conditionsefor integra\$Qrs


## Examples of first integrals

- Conservation of energy in Hamiltonian systems


## Hamiltonian system

$$
\dot{p}=-\frac{\partial H}{\partial q}, \quad \dot{q}=\frac{\partial H}{\partial p}
$$

## Theorem

$$
\frac{d}{d t} H(p, q)=\frac{\partial H}{\partial p} \dot{p}+\frac{\partial H}{\partial q} \dot{q}=0 \text { hence } H(p, q)=\text { Const }
$$

## General invariants encountered in physics

## Examples of first integrals

- Conservation of energy in Hamiltonian systems
- Conservation of total and angular momentum in N-Body systems


## N-Body system

$$
\dot{p}_{i}=-\sum_{j=1}^{N} \nu_{i j}\left(q_{i}-q_{j}\right), \quad \dot{q}_{i}=\frac{p_{i}}{m_{i}} \quad \nu \text { symmetric }
$$

## Theorem

$\sum_{i=1}^{N} p_{i}=$ Const and $\sum_{i=1}^{N} q_{i} \times p_{i}=$ Const

## General invariants encountered in physics

## Examples of first integrals

- Conservation of energy in Hamiltonian systems
- Conservation of total and angular momentum in N-Body systems
- Conservation of mass in chemical reactions


## Chemical reactions

$$
\begin{array}{cccccc}
A & \xrightarrow{k_{1}} & B & \dot{y}_{1}= & -k_{1} y_{1}+k_{3} y_{2} y_{3} \\
B+B & \xrightarrow{k_{2}} & B+C & \dot{y}_{2} & = & k_{1} y_{1}-k_{3} y_{2} y_{3}-k_{2} y_{2}^{2} \\
B+C & \xrightarrow{k_{3}} & A+C & \dot{y}_{3} & = & k_{2} y_{2}^{2}
\end{array}
$$

## Theorem

$$
\frac{d}{d t}\left(y_{1}+y_{2}+y_{3}\right)=0 \text { hence } I(y)=y_{1}+y_{2}+y_{2}=\text { Const. }
$$

## General invariants encountered in physics

## Examples of first integrals

- Conservation of energy in Hamiltonian systems
- Conservation of total and angular momentum in N-Body systems
- Conservation of mass in chemical reactions
- Conservation of the spectrum by matrix flows


## Isospectral matrix equations

$$
\dot{L}=B(L) L-L B(L) \text { with } B(L) \text { skew-symmetric. }
$$

## Theorem

Let $\dot{U}=B(L(t)) U, U(0)=I$. Then, $L(t)=U(t) L_{0} U(t)^{-1}$.

## General invariants encountered in physics

## Examples of first integrals

- Conservation of energy in Hamiltonian systems
- Conservation of total and angular momentum in N-Body systems
- Conservation of mass in chemical reactions
- Conservation of the spectrum by matrix flows
- Conservation of volume in divergence-free systems


## Divergence-free system

$$
\dot{y}=f(y) \text { with } \operatorname{div}(f)=0 \text {. }
$$

## Theorem

The flow $\varphi_{t}$ preserves the volume, i.e. $\int_{\varphi_{t}(A)} d y=\int_{A} d y$.

## A prey-predator model in normal form

$$
\begin{aligned}
\dot{U} & =e^{V}-2=f(V) \\
\dot{V} & =1-e^{U}=g(U)
\end{aligned}
$$



## A prey-predator model in normal form

$$
\begin{aligned}
\dot{U} & =e^{V}-2=f(V) \\
\dot{V} & =1-e^{U}=g(U)
\end{aligned}
$$



## 2-D Kepler Problem

$$
H(p, q)=\frac{1}{2} p^{T} p-\frac{1}{\sqrt{q^{T} q}}=T(p)+V(q) \Longleftrightarrow \ddot{q}=-V^{\prime}(q)
$$

Euler explicit/implicit


## 2-D Kepler Problem

$$
H(p, q)=\frac{1}{2} p^{T} p-\frac{1}{\sqrt{q^{T} q}}=T(p)+V(q) \Longleftrightarrow \ddot{q}=-V^{\prime}(q) .
$$



## 2-D Kepler Problem

$$
H(p, q)=\frac{1}{2} p^{T} p-\frac{1}{\sqrt{q^{T} q}}=T(p)+V(q) \Longleftrightarrow \ddot{q}=-V^{\prime}(q) .
$$



## Outline

(1)

## Problems and motivations

- General invariants encountered in physics
- Improved qualitative behavior of geometric integrators


## (2) Setting of the problem

- Invariant and volume preservation for split systems
- The Hopf algebra of coloured trees
- B-series and S-series for split vector fields

Conditions for invariants-preservation

- Numerical methods preserving invariants
- The case of quadratic and cubic invariants
- B-series methods preserving all cubic invariants

Conditions for volume-preservation

- Volume-preserving B-series
- Connection with the preservation of cubic invariants
- volume preserving methods for split systems with a special stru From conditions for vector fields to conditionsefor integra\$Qrs


## The two classes of problems considered

We consider systems of ODEs of the form

## Split vector fields systems

$$
\dot{y}=f^{[1]}(y)+f^{[2]}(y)+\ldots+f^{[N]}(y),
$$

such that each individual vector field has the invariant function /
Common Invariant

$$
0=\left(\nabla_{y} l(y)\right)^{\top} f^{[\nu]}(y), \quad \nu=1, \ldots, N,
$$

## The two classes of problems considered

We consider systems of ODEs of the form

## Split vector fields systems

$$
\dot{y}=f^{[1]}(y)+f^{[2]}(y)+\ldots+f^{[N]}(y),
$$

or preserves the volume form

## Divergence-free

$$
0=\operatorname{div} f^{[\nu]}(y), \quad \nu=1, \ldots, N
$$

## Invariant preserving integrators

## A one-step method

is a map from the phase-space to itself, which, given an approximation $y$ of the solution at time $t$, produces an approximation $\Phi_{h}^{f}$ at time $t+h$.

## The modified vector field

 associated to a numerical integrator $\Phi_{h}^{f}$ is the vector field $\tilde{f}_{h}$ such that the exact solution of $\dot{z}=\tilde{f}_{h}(z), z(t)=y$ at time $t+h$ is $\Phi_{h}^{f}(y)$.
## Invariant-preserving integrators(1)

$\Phi_{h}^{f}$ preserves $I$ if $I\left(\Phi_{h}^{f}(y)\right)=I(y)$ for any $y$.

## Invariant preserving integrators

## A one-step method

is a map from the phase-space to itself, which, given an approximation $y$ of the solution at time $t$, produces an approximation $\Phi_{h}^{f}$ at time $t+h$.

## The modified vector field

 associated to a numerical integrator $\Phi_{h}^{f}$ is the vector field $\tilde{f}_{h}$ such that the exact solution of $\dot{z}=\tilde{f}_{h}(z), z(t)=y$ at time $t+h$ is $\Phi_{h}^{f}(y)$.
## Invariant-preserving integrators(2)

$\Phi_{h}^{f}$ preserves $/$ if $(\nabla /(y))^{T} \tilde{f}_{h}(y)=0$ for any $y$.

## Volume-preserving integrators

## A one-step method

is a map from the phase-space to itself, which, given an approximation $y$ of the solution at time $t$, produces an approximation $\Phi_{h}^{f}(y)$ at time $t+h$.

## The modified vector field

 associated to a numerical integrator $\Phi_{h}^{f}$ is the vector field $\tilde{f}_{h}$ such that the exact solution of $\dot{z}=\tilde{f}_{h}(z), z(t)=y$ at time $t+h$ is $\Phi_{h}^{f}(y)$.
## Volume-preserving integrators(1)

$\Phi_{h}^{f}$ preserves the volume if $\operatorname{det}\left(\frac{\partial \Phi_{h}^{f}(y)}{\partial y}\right)=1$ for any $y$.

## Volume-preserving integrators

## A one-step method

is a map from the phase-space to itself, which, given an approximation $y$ of the solution at time $t$, produces an approximation $\Phi_{h}^{f}(y)$ at time $t+h$.

## The modified vector field

 associated to a numerical integrator $\Phi_{h}^{f}$ is the vector field $\tilde{f}_{h}$ such that the exact solution of $\dot{z}=\tilde{f}_{h}(z), z(t)=y$ at time $t+h$ is $\Phi_{h}^{f}(y)$.
## Volume-preserving integrators(2)

$\Phi_{h}^{f}$ preserves the volume if $\operatorname{div}\left(\tilde{f}_{h}(y)\right)=0$ for any $y$.

## Volume-preserving integrators

## A one-step method

is a map from the phase-space to itself, which, given an approximation $y$ of the solution at time $t$, produces an approximation $\Phi_{h}^{f}(y)$ at time $t+h$.

## The modified vector field

 associated to a numerical integrator $\Phi_{h}^{f}$ is the vector field $\tilde{f}_{h}$ such that the exact solution of $\dot{z}=\tilde{f}_{h}(z), z(t)=y$ at time $t+h$ is $\Phi_{h}^{f}(y)$.The conditions for preserving the volume are easier to obtain in terms of the modified vector field.

## The Hopf algebra of coloured trees

## Trees and forests [Merson 57, Butcher 68]

## Definition

The set of trees $\mathcal{T}$ and forests $\mathcal{F}$ are defined recursively by:
(1) $e \in \mathcal{F}$
(2) if $t_{1}, \ldots, t_{n} \in \mathcal{T}^{n}$, then $u=t_{1} \ldots t_{n} \in \mathcal{F}$
(3) if $u \in \mathcal{F}$ and $\nu \in\{1, \ldots, N\}$, then $t=[u]_{\nu}=B_{\nu}^{+}(u) \in \mathcal{T}$.

## Example

$$
\begin{array}{r}
B_{1}^{+}(\cdot \circ)=[\cdot \circ]_{1}=\vartheta \text { and } B_{2}^{+}(\cdot \cdot)=[\cdots]_{2}=\vartheta \\
B^{-}(\mathscr{V})=\cdots \text { and } B^{-}(\bigvee)=\vartheta
\end{array}
$$

## The Hopf algebra of coloured trees

## Order and symmetry

## Definition

Consider $n$ distinct trees $t_{1}, \ldots, t_{n}$ and let $u=t_{1}^{t_{1}} \ldots t_{n}^{r_{n}}$ and $t=[u]_{\nu}$. Then,

- $|t|=1+|u|=1+r_{1}\left|t_{1}\right|+\ldots+r_{n}\left|t_{n}\right|$
- $\sigma(u)=r_{1}!\ldots r_{n}!\left(\sigma\left(t_{1}\right)\right)^{r_{1}} \ldots\left(\sigma\left(t_{n}\right)\right)^{r_{n}}$ and $\sigma(t)=\sigma(u)$


## Example

| Forest $u$ | -•! | $v \vee!$ | $\%^{33} \mathrm{~V}$ 彦 | v:8 |
| :---: | :---: | :---: | :---: | :---: |
| Order \|u| | 4 | 11 | 17 | 11 |
| Symmetry $\sigma(u)$ | $2!$ | 1!3!1! | $3!(2!)^{3} 2$ ! | $3!1!1!$ |

## The Hopf algebra of coloured trees

## Structure (Connes and Kreimer 98, Brouder 04)

## Definition

The set $\mathcal{F}$ can be naturally endowed with an algebra structure $\mathcal{H}$ on $\mathbb{R}$ :

- $\forall(u, v) \in \mathcal{F}^{2}, \forall(\lambda, \mu) \in \mathbb{R}^{2}, \lambda u+\mu v \in \mathcal{H}$,
- $\forall(u, v) \in \mathcal{F}^{2}, u v \in \mathcal{H}$ (note that $u v=v u$ ),
- $\forall u \in \mathcal{F}, u e=e u=u$.


## Calculus in $\mathcal{H}$



## The Hopf algebra of coloured trees

## The co-product

## Definition

The tensor product of $\mathcal{H}$ with itself is the set of elements of the form $u \otimes v$ such that for all $(u, v, w, x) \in \mathcal{H}^{4}$ and all $(\lambda, \mu) \in \mathbb{R}^{2}$ :

$$
\begin{aligned}
& (\lambda u+\mu v) \otimes w=\lambda(u \otimes w)+\mu(v \otimes w), \\
& w \otimes(\lambda u+\mu v)=\lambda(w \otimes u)+\mu(w \otimes v), \\
& (u \otimes v)(w \otimes x)=(u w \otimes v x) .
\end{aligned}
$$

## Definition

The co-product $\Delta$ is a morphism from $\mathcal{H}$ to $\mathcal{H} \otimes \mathcal{H}$ defined by:
(1) $\Delta(e)=e \otimes e$,
(2) $\forall t \in \mathcal{T}, \Delta(t)=t \otimes e+\left(i d_{\mathcal{H}} \otimes B_{\mu(t)}^{+}\right) \circ \Delta \circ B^{-}(t)$,
(3) $\forall u=t_{1} \ldots t_{n} \in \mathcal{F}, \Delta(u)=\Delta\left(t_{1}\right) \ldots \Delta\left(t_{n}\right)$.

## The Hopf algebra of coloured trees

## The co-product

## Example

$$
\begin{aligned}
& \Delta\left(\gamma^{\circ}\right)=\quad{ }^{\circ} \otimes e+\left(i d \otimes B_{2}^{+}\right) \Delta(\cdot \circ) \\
& =\dot{\gamma} \otimes e+\left(i d \otimes B_{2}^{+}\right) \Delta(\cdot) \Delta(\circ) \\
& =\vartheta \bullet \otimes e+\left(i d \otimes B_{2}^{+}\right)(\cdot \otimes e+e \otimes \cdot)(\circ \otimes e+e \otimes \circ) \\
& =\dot{\gamma} \otimes e \\
& +\left(i d \otimes B_{2}^{+}\right)(\cdot \circ \otimes e+\cdot \otimes \circ+\circ \otimes \cdot+\boldsymbol{e} \otimes \cdot \circ)
\end{aligned}
$$

## B-series and S-series for split vector fields

## Elementary differentials

## Definition

Let $t$ be a tree of $\mathcal{T}$. The elementary differential $F(t)$ associated with $t$ is the mapping from $\mathbb{R}^{n}$ to $\mathbb{R}^{n}$, defined by:
(1) $F\left(\cdot{ }_{\nu}\right)(y)=f^{[\nu]}(y)$,
(2) $F\left(\left[t_{1}, \ldots, t_{n}\right]_{\nu}\right)(y)=\left(f^{[\nu]}\right)^{(n)}(y)\left(F\left(t_{1}\right)(y), \ldots, F\left(t_{n}\right)(y)\right)$.

## Example

$$
\begin{aligned}
F(\dot{\ell}) & =\left(f^{[1]}\right)^{\prime} f^{[2]} \\
F\left(\mathfrak{\zeta}^{\circ}\right) & =\left(f^{[2]}\right)^{\prime \prime}(y)\left(f^{[1]}, f^{[2]}\right) \\
F(\dot{\varrho}) & =\left(f^{[1]}\right)^{\prime}\left(f^{[2]}\right)^{\prime} f^{[1]}
\end{aligned}
$$

## B-series and S-series for split vector fields

## Elementary differential operators

## Definition

Let $u=t_{1} \ldots t_{k}$ be a forest of $\mathcal{F}$. The differential operator $X(u)$ associated with $u$ is defined on $\mathcal{D}=C^{\infty}\left(\mathbb{R}^{n} ; \mathbb{R}^{m}\right)$ by:

$$
\begin{aligned}
X(u): \mathcal{D} & \rightarrow \mathcal{D} \\
g & \mapsto X(u)[g]=g^{(k)}\left(F\left(t_{1}\right), \ldots, F\left(t_{k}\right)\right) .
\end{aligned}
$$

## Example

$$
\begin{aligned}
X(e)[g] & =g \\
X(\cdot)[g] & =g^{\prime} f^{[1]} \\
X(f)[g] & =g^{\prime}\left(f^{[1]}\right)^{\prime} f^{[2]} \\
X(f \circ \cdot)[g] & =g^{(3)}\left(\left(f^{[1]}\right)^{\prime} f^{[1]}, f^{[2]}, f^{[1]}\right)
\end{aligned}
$$

## B-series and S-series for split vector fields

## B-series and S-series

## Definition (B-Series (Hairer and Wanner 74))

Let $a: \mathcal{T} \rightarrow \mathbb{R}$. The B -series $B(a, y)$ is the formal series:

$$
B(a, y)=a(e) y+\sum_{t \in \mathcal{T}} \frac{h^{|t|}}{\sigma(t)} a(t) F(t)
$$

## Example (Implicit/Explicit Euler)

$$
\begin{aligned}
y_{1}= & y_{0}+h\left(f^{[1]}\left(y_{1}\right)+f^{[2]}\left(y_{0}\right)\right) \\
= & y_{0}+h F(\cdot)\left(y_{0}\right)+h F(\circ)\left(y_{0}\right)+h^{2} F(\cdot)\left(y_{0}\right)+h^{2} F(\rho)\left(y_{0}\right) \\
& +\ldots
\end{aligned}
$$

## B-series and S-series for split vector fields

## B-series and S-series

## Definition (Series of differential operators)

Let $\alpha: \mathcal{F} \rightarrow \mathbb{R}$. The S-series $S(\alpha)$ is the formal series

$$
S(\alpha)[g]=\sum_{u \in \mathcal{F}} \frac{h^{|u|}}{\sigma(u)} \alpha(u) X(u)[g]
$$

## Example (Implicit/Explicit Euler)

$$
\begin{aligned}
g\left(y_{1}\right)= & g\left(y_{0}+h f^{[1]}\left(y_{1}\right)+h f^{[2]}\left(y_{0}\right)\right) \\
= & X(e)[g]+h(X(\cdot)[g]+X(\circ)[g])+h^{2}(X(!)[g]+X(\circ)[g]) \\
& +\frac{h^{2}}{2}\left(X\left(\cdot^{2}\right)[g]+2 X(\cdot \circ)[g]+X\left(\circ^{2}\right)[g]\right)+\ldots
\end{aligned}
$$

## B-series and S-series for split vector fields

## Composition of series and co-product in $\mathcal{H}$

## Theorem (Composition of B-series)

Let $a$ and $b$ be two mappings from $\mathcal{T}$ to $\mathbb{R}$. The composition of the two $B$-series $B(a, y)$ and $B(b, y)$, i.e. $B(b, B(a, y))$, is again a $B$-series $B(a . b, y)$, with coefficients a.b defined on $\mathcal{T}$ by

$$
\forall t \in \mathcal{T}, \quad(a . b)(t)=\left(\mu_{\mathbb{R}} \circ(a \otimes b) \circ \Delta\right)(t) .
$$

## Example

$$
\begin{aligned}
& (a . b)(\mathcal{O})=\mu_{\mathbb{R}} \circ(a \otimes b)(\hat{O} \otimes e+\cdots \otimes 0+\bullet \otimes \boldsymbol{O}+o \otimes \sigma+e \otimes \dot{O}) \\
& =a(\mathrm{O}) b(e)+a(\bullet) a(\circ) b(\circ)+a(\bullet) b\left(\delta_{0}\right)+a(\circ) b(\delta)+a(e) b(\text { К })
\end{aligned}
$$

## B-series and S-series for split vector fields

## Composition of series and co-product in $\mathcal{H}$

## Theorem (Composition of S-series)

Let $\alpha$ and $\beta$ be two mappings from $\mathcal{F}$ to $\mathbb{R}$. The composition of the two $S$-series $S(\alpha)$ and $S(\beta)$, i.e. $S(\alpha)[S(\beta)[]$.$] is again a$ $S$-series, with coefficients $\alpha . \beta$ defined on $\mathcal{F}$ by

$$
\forall u \in \mathcal{F}, \quad(\alpha \beta)(u)=\left(\mu_{\mathbb{R}} \circ(\alpha \otimes \beta) \circ \Delta\right)(u)
$$

## Example

$$
\begin{aligned}
& (\alpha . \beta)(\bigvee)=\mu_{\mathbb{R}} \circ(\alpha \otimes \beta)(\bigvee \otimes e+\cdots \otimes \circ+\cdots \otimes \boldsymbol{O}+\circ \otimes \boldsymbol{\sigma}+e \otimes \hat{\zeta})
\end{aligned}
$$

## Outline

```
Problems and motivations
- General invariants encountered in physics
- Improved qualitative behavior of geometric integrators
Setting of the problem
- Invariant and volume preservation for split systems
- The Hopf algebra of coloured trees
- B-series and S-series for split vector fields
```

(3) Conditions for invariants-preservation

- Numerical methods preserving invariants
- The case of quadratic and cubic invariants
- B-series methods preserving all cubic invariants

Conditions for volume-preservation

- Volume-preserving B-series
- Connection with the preservation of cubic invariants
- volume preserving methods for split systems with a special stru From conditions for vector fields to conditionsefor integrawQrs


## The action of a function / on a B-series

It can be viewed as S-series:

$$
I(B(a, y))=S(\alpha)\left[I \Longleftrightarrow \alpha \in \operatorname{Alg}(\mathcal{H}, \mathbb{R}) \text { and } \alpha_{\mid \mathcal{T}} \equiv a .\right.
$$

## A B-series integrator $B(a, y)$ preserves $/$ iff

$$
\forall y \in \mathbb{R}^{n}, I(B(a, y))=I(y),
$$

i.e.

$$
S(\alpha)[I]=I,
$$

where $\alpha$ is the unique algebra-morphism extending $a$ onto $\mathcal{H}$.

## The annihilating left ideal $\mathcal{I}[/]$ of /

## Using the assumption of a common invariant /

For $\nu=1, \ldots, N, X\left(\cdot{ }_{\nu}\right)[l]=(\nabla I) f^{[\nu]}=0$. Hence,

$$
\sum^{N} S\left(\omega_{\nu}\right)\left[h X\left(\cdot{ }_{\nu}\right)[I]\right]=S\left(\omega^{\prime}\right)[I]=0 .
$$

## Lemma

For any $\left(\omega_{1}, \ldots, \omega_{N}\right) \in\left(\mathcal{H}^{*}\right)^{N}$, we have $\omega^{\prime}(e)=0$ and

$$
\forall u=t_{1} \cdots t_{m} \in \mathcal{F}, \quad \omega^{\prime}(u)=\sum_{i=1}^{m} \omega_{\mu\left(t_{i}\right)}\left(B^{-}\left(t_{i}\right) \prod_{j \neq i} t_{j}\right) .
$$

## Numerical methods preserving invariants

## Integrators preserving general invariants

## Theorem

Let $\alpha \in \operatorname{Alg}(\mathcal{H}, \mathbb{R})$. Then $\alpha$ satisfies $S(\alpha)[I]=I$ that for all couples $(f, I)$ of a vector field $f$ and a first integral I, if and only if $\alpha(e)=1$ and $\alpha$ satisfies the condition

$$
\alpha\left(t_{1}\right) \cdots \alpha\left(t_{m}\right)=\sum_{j=1}^{m} \alpha\left(t_{j} \circ \prod_{i \neq j} t_{i}\right)
$$

for all $m \geq 2$ and all $t_{i}$ 's in $\mathcal{T}$.

## Theorem

Let $\beta \in \operatorname{VF}(\mathcal{H}, \mathbb{R})$. Then $\beta$ satisfies $S(\beta)[I]=0$ that for all couples $(f, I)$ if and only if $\alpha$ satisfies the condition

$$
0=\sum_{j=1}^{m} \beta\left(t_{j} \circ \prod_{i \neq j} t_{i}\right)
$$

## For quadratic

first integral $I$, the condition becomes

$$
\forall\left(t_{1}, t_{2}\right) \in \mathcal{T}^{2}, \quad b\left(t_{1} \circ t_{2}\right)+b\left(t_{2} \circ t_{1}\right)=0
$$

while for cubic invariants $I$, one needs in addition that

$$
\forall\left(t_{1}, t_{2}, t_{3}\right) \in \mathcal{T}^{3}, \quad b\left(t_{1} \circ t_{2} t_{3}\right)+b\left(t_{2} \circ t_{1} t_{3}\right)+b\left(t_{3} \circ t_{1} t_{2}\right)=0
$$

## Theorem

A B-series integrator that preserves all cubic polynomial invariants does in fact preserve polynomial invariants of any degree and can be formally interpreted as the exact flow of a vector field lying in the Lie-algebra generated by $f^{[1]}, \ldots, f^{[N]}$.

## Outline

(1)

## Problems and motivations

- General invariants encountered in physics
- Improved qualitative behavior of geometric integrators Setting of the problem
- Invariant and volume preservation for split systems
- The Hopf algebra of coloured trees
- B-series and S-series for split vector fields

Conditions for invariants-preservation

- Numerical methods preserving invariants
- The case of quadratic and cubic invariants
- B-series methods preserving all cubic invariants

4 Conditions for volume-preservation

- Volume-preserving B-series
- Connection with the preservation of cubic invariants
- volume preserving methods for split systems with a special stru


## Split systems with zero-divergence

## Divergence-free B-series

For systems of the form

$$
\dot{y}=\sum_{\nu=1}^{N} f^{[\nu]}(y) \text { with } \operatorname{div} f^{[\nu]}=0
$$

a B-series modified vector field is divergence free if

$$
\operatorname{div}\left(h \tilde{f}_{h}(y)\right)=\sum_{t \in \mathcal{T}} \frac{h^{|t|}}{\sigma(t)} b(t) \operatorname{div}(F(t)(y))=0
$$

## Question

How to compute and represent the terms in $\operatorname{div}(F(t)(y) ?$

## Volume-preserving B-series

## A convenient formula for the derivative of an elementary differential

## Notation

For $t=\left[t_{1}, \ldots, t_{l}\right]_{\nu} \in \mathcal{T}, \quad F^{*}(t)=\frac{\partial^{I+1} f^{[\nu]}}{\partial y^{\prime+1}}\left(F\left(t_{1}\right), \ldots, F\left(t_{l}\right)\right)$.

## The formula

$$
\frac{d F(t)}{\sigma(t)}=\frac{F^{*}(t)}{\sigma(t)}+\sum_{t_{1} \hat{t}_{2}, \ldots o t_{m}=t} \frac{F^{*}\left(t_{1}\right)}{\sigma\left(t_{1}\right)} \frac{F^{*}\left(t_{2}\right)}{\sigma\left(t_{2}\right)} \cdots \frac{F^{*}\left(t_{m}\right)}{\sigma\left(t_{m}\right)} .
$$

The grafting operation is meant to operate from right to left.

$$
\frac{\operatorname{div}(F(t)}{\sigma(t)}=\frac{\operatorname{Tr}\left(F^{*}(t)\right)}{\sigma(t)}+\sum_{t_{1} \circ t_{2} \circ \cdots \circ t_{m}=t} \frac{\operatorname{Tr}\left(F^{*}\left(t_{1}\right) \ldots F^{*}\left(t_{m}\right)\right)}{\sigma\left(t_{1}\right) \ldots \sigma\left(t_{m}\right)}
$$

## The set of aromatic trees $\mathcal{A T}$

## Definition

An aromatic tree $o$ is a coloured oriented graph with exactly one cycle, such that if all the arcs in the cycle are removed, then the resulting coloured oriented graph is identified with a forest $t_{1} \cdots t_{m}$. If the arcs of $o$ that form the cycle go from the root of $t_{i}$ to the root of $t_{i+1}(i=1, \ldots, m-1)$ and from the root of $t_{m}$ to the root of $t_{1}$ then we write $o=\left(t_{1} \cdots t_{m}\right)$. The set of aromatic trees is denoted $\mathcal{A T}$ and the set of $n$-th order aromatic trees $\mathcal{A} \mathcal{T}_{n}$.
$0=\underset{0 \rightarrow 0}{\substack{0 \\ 0 \rightarrow 0}}=\left(t_{1} t_{2} t_{1} t_{2}\right) \quad t_{1}=0 \rightarrow 0<0=$.

## 1-cuts of aromatic trees

## Definition

For any aromatic tree $o=\left(t_{1} \ldots t_{m}\right) \in \mathcal{A T}, C(o)$ is the unordered list of trees obtained from o by breaking any edge of the cycle. If we denote for $i=1, \ldots, m$, $s_{i}=t_{i} \circ t_{i+1} \circ \ldots \circ t_{m} \circ t_{1} \circ \ldots \circ t_{i-1}$, then:

$$
\begin{equation*}
C(o)=\left\{s_{1}, \ldots, s_{m}\right\} . \tag{1}
\end{equation*}
$$

Now, let $\pi_{m}$ be the circular permutation of $\{1, \ldots, m\}$ and let $\theta$ be
$\theta=\#\left\{I \in\{0, \ldots, m-1\}: \quad\left(t_{\pi_{m}^{\prime}(1)}, \ldots, t_{\pi_{m}^{\prime}(m)}\right)=\left(t_{1}, \ldots, t_{m}\right)\right\}$,
so that, for each $i$, there are $\theta$ copies of $s_{i}$ in the list $C(o)$. Then the symmetry coefficient of $o$ is defined as $\sigma(o)=\theta \prod_{i} \sigma\left(t_{i}\right)$.

## Volume-preserving B-series

## The list $C(0)=\left\{s_{1}, s_{2}, s_{3}, s_{4}\right\}$ for $0=\left(t_{1} t_{2} t_{1} t_{2}\right)$



## Divergence of a B-series vector field

## Definition (Elementary divergence)

The divergence div(o) associated with an aromatic tree $o=\left(t_{1} \ldots t_{m}\right)$ is defined by:

$$
\operatorname{div}(0)=\operatorname{Tr}\left(F^{*}\left(t_{1}\right) \ldots F^{*}\left(t_{m}\right)\right) .
$$

## Collecting the terms

$$
\begin{aligned}
\operatorname{div}(B(b)) & =\sum_{t \in \mathcal{T}} b(t) h^{|t|} \sum_{m \geq 2} \sum_{t_{1} \cdots \cdots t_{m}=t} \frac{\operatorname{div}\left(\left(t_{1} \ldots t_{m}\right)\right)}{\sigma\left(t_{1}\right) \cdots \sigma\left(t_{m}\right)} \\
& =\sum_{n \geq 2} h^{n} \sum_{o \in \mathcal{A} \mathcal{T}_{n}}\left(\sum_{t \in C(o)} b(t)\right) \frac{\operatorname{div}(o)}{\sigma(o)} .
\end{aligned}
$$

## Divergence-free conditions

## Theorem

A modified field given by the $B$-series $B(b, y)$ is divergence-free up to order $p$ if the following condition is satisfied:

$$
\sum_{t \in C(o)} b(t)=0 \text { for all } o \in \mathcal{A T} \text { with }|o| \leq p
$$

## Example

For $0=\left(t_{1} t_{2} t_{1} t_{2}\right)$,

$$
2 b\left(t_{1} \circ t_{2} \circ t_{1} \circ t_{2}\right)+2 b\left(t_{2} \circ t_{1} \circ t_{2} \circ t_{1}\right)=0
$$

## 2-3 cycles conditions and conditions for quadratic/cubic invariants

(1) 2-cycles clearly coincide with the conditions for quadratic invariants.
(2) for 3-cycles conditions

$$
\begin{aligned}
0 & =b\left(t_{1} \circ t_{2} \circ t_{3}\right)+b\left(t_{2} \circ t_{1} \circ t_{3}\right)+b\left(t_{3} \circ t_{2} \circ t_{1}\right), \\
& =b\left(t_{1} \circ\left(t_{2} \circ t_{3}\right)\right)+b\left(t_{2} \circ\left(t_{1} \circ t_{3}\right)\right)+b\left(t_{3} \circ\left(t_{2} \circ t_{1}\right)\right), \\
& =-b\left(\left(t_{2} \circ t_{3}\right) \circ t_{1}\right)-b\left(\left(t_{1} \circ t_{3}\right) \circ t_{2}\right)-b\left(\left(t_{2} \circ t_{1}\right) \circ t_{3}\right), \\
& =-b\left(t_{2} \circ t_{1} t_{3}\right)-b\left(t_{1} \circ t_{2} t_{3}\right)-b\left(t_{2} \circ t_{1} t_{3}\right) .
\end{aligned}
$$

## Theorem

A volume-preserving B-series integrator can be formally interpreted as the exact flow of a vector field lying in the Lie-algebra generated by $f^{[1]}, \ldots, f^{[N]}$.

## The conditions for a special class of systems

## 3-cycle systems

$$
\left(\begin{array}{l}
\dot{p} \\
\dot{q} \\
\dot{r}
\end{array}\right)=\left(\begin{array}{c}
\mathcal{F}(q) \\
\mathcal{G}(r) \\
\mathcal{H}(p)
\end{array}\right)=f^{[1]}(q)+f^{[2]}(r)+f^{[3]}(p) .
$$

## Black trees

For $u=\left[v_{1}, \ldots, v_{m}\right]$. , one has

$$
F^{*}(u)=\frac{\partial^{m+1} f^{[1]}}{\partial(p, q, r)^{m+1}}\left(F\left(v_{1}\right), \ldots, F\left(v_{m}\right)\right)=\left(\begin{array}{ccc}
0 & \times & 0 \\
0 & 0 & 0 \\
0 & 0 & 0
\end{array}\right)
$$

## The conditions for a special class of systems

## 3-cycle systems

$$
\left(\begin{array}{c}
\dot{p} \\
\dot{q} \\
\dot{r}
\end{array}\right)=\left(\begin{array}{c}
\mathcal{F}(q) \\
\mathcal{G}(r) \\
\mathcal{H}(p)
\end{array}\right)=f^{[1]}(q)+f^{[2]}(r)+f^{[3]}(p) .
$$

## White trees

For $v=\left[w_{1}, \ldots, w_{n}\right]_{\circ}$, one has

$$
F^{*}(v)=\frac{\partial^{n+1} f^{[2]}}{\partial(p, q, r)^{n+1}}\left(F\left(w_{1}\right), \ldots, F\left(w_{n}\right)\right)=\left(\begin{array}{ccc}
0 & 0 & 0 \\
0 & 0 & \times \\
0 & 0 & 0
\end{array}\right) .
$$

## The conditions for a special class of systems

## 3-cycle systems

$$
\left(\begin{array}{c}
\dot{p} \\
\dot{q} \\
\dot{r}
\end{array}\right)=\left(\begin{array}{c}
\mathcal{F}(q) \\
\mathcal{G}(r) \\
\mathcal{H}(p)
\end{array}\right)=f^{[1]}(q)+f^{[2]}(r)+f^{[3]}(p) .
$$

## Square trees

For $w=\left[u_{1}, \ldots, u_{r}\right] \square$, one has

$$
F^{*}(w)=\frac{\partial^{r+1} f^{[3]}}{\partial(p, q, r)^{r+1}}\left(F\left(w_{1}\right), \ldots, F\left(w_{n}\right)\right)=\left(\begin{array}{ccc}
0 & 0 & 0 \\
0 & 0 & 0 \\
\times & 0 & 0
\end{array}\right) .
$$

## The conditions for a special class of systems

## 3-cycle systems

$$
\left(\begin{array}{l}
\dot{p} \\
\dot{q} \\
\dot{r}
\end{array}\right)=\left(\begin{array}{c}
\mathcal{F}(q) \\
\mathcal{G}(r) \\
\mathcal{H}(p)
\end{array}\right)=f^{[1]}(q)+f^{[2]}(r)+f^{[3]}(p) .
$$

## Consequence

$$
\operatorname{div}(o) \neq 0 \text { iff } o=\left(u_{1} v_{1} w_{1} u_{2} v_{2} w_{2} \ldots u_{m} v_{m} w_{m}\right), m \geq 1 .
$$

## Volume-preserving RK-methods for 3-cycle systems

## Theorem

A one-stage additive Runge-Kutta method formed of $\left(A^{[]]}, b^{[l]}\right)=\left(\theta_{i}, 1\right), i=1,2,3$, is volume-preserving for 3 -cycle systems iff

$$
\left(\theta_{1}-1\right)\left(\theta_{2}-1\right)\left(\theta_{3}-1\right)=\theta_{1} \theta_{2} \theta_{3}
$$

## Example

An implicit "non-symplectic" RK-method

$$
\begin{aligned}
P & =p_{0}+\frac{h}{3} \mathcal{F}(Q) \quad p_{1}=p_{0}+h \mathcal{F}(Q) \\
Q & =q_{0}+\frac{4 h}{3} \mathcal{G}(R) \quad q_{1}=q_{0}+h \mathcal{G}(R) \\
R & =r_{0}+\frac{h}{3} \mathcal{H}(P) \quad p_{1}=p_{0}+h \mathcal{H}(P)
\end{aligned}
$$

## Outline

1

## Problems and motivations

- General invariants encountered in physics
- Improved qualitative behavior of geometric integrators Setting of the problem
- Invariant and volume preservation for split systems
- The Hopf algebra of coloured trees
- B-series and S-series for split vector fields
(3) Conditions for invariants-preservation
- Numerical methods preserving invariants
- The case of quadratic and cubic invariants
- B-series methods preserving all cubic invariants
(4) Conditions for volume-preservation
- Volume-preserving B-series
- Connection with the preservation of cubic invariants
- volume preserving methods for split systems with a special stru
(5) From conditions for vector fields to conditions for integrators


## Substitution law

## From integrators to vector fields and vice-versa

Backward error analysis


## Back to the black forest

Though what follows is valid for multicoloured trees, for simplicity we now turn back to the monocolour situation.

## Substitution law

## From partitions and skeletons to the formula

## Definition

Given a partition $p$ of $t$, the corresponding skeleton $\chi_{p}$ is the tree obtained by contracting each tree of $p$ to a single vertex and by re-establishing the cut edges.

Table: The 8 partitions of a tree of order 4 with associated skeleton and forest

| $p$ | $\ddots$ | $\ddots$ | $\ddots$ | $\ddots$ | $\ddots$ | $\ddots$ | $\ddots$ | $\ddots$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $x_{p}$ | $\cdot$ | $\ddots$ | $\ddots$ | $\ddots$ | $\ddots$ | $\ddots$ | $\vdots$ | $\ddots$ |
| $v_{p}$ | $\ddots$ | $\cdot \vartheta$ | $\cdot \ddots$ | $\cdot \ddots$ | $\cdots$ | $\cdots!$ | $\cdots!$ | $\cdots$ |

## Theorem

For $b(\emptyset)=0$, the vector field $h^{-1} B_{f}(b, y)$ inserted into $B_{g}(a, y)$, i.e. with $g=h^{-1} B_{f}(b, y)$ gives a $B$-series

$$
B_{g}(a, y)=B_{f}(b \star a, y)
$$

We have $(b \star a)(\emptyset)=a(\emptyset)$ and for all $t \in \mathcal{T}$,

$$
(b \star a)(t)=\sum_{p \in \mathcal{P}(t)} a\left(\chi_{p}\right) b\left(v_{p}\right)
$$

Table: Substitution law $\star$ for the first trees.

$$
\begin{aligned}
(b \star a)(\emptyset) & =a(\emptyset) \\
(b \star a)(\cdot) & =a(\cdot) b(\cdot) \\
(b \star a)(!) & =a(\cdot) b(!)+a(!) b(\cdot)^{2} \\
(b \star a)(\zeta) & =a(\cdot) b(\zeta)+2 a(!) b(\cdot) b(!)+a(\zeta) b(\cdot)^{3} \\
(b \star a)(!) & =a(\cdot) b(!)+2 a(!) b(\cdot) b(!)+a(!) b(\cdot)^{3}
\end{aligned}
$$

## Remark

This law essentially coincides with the convolution product in the Hopf algebra of Calaque, Ebrahimi-Fard and Manchon.

Let $\omega$ denote the inverse element of $\frac{1}{\gamma}-\delta_{\emptyset}$ for $\star$. The backward error coefficients $b$ can be computed as follows:

## Backward error character $\omega$

$$
\forall t \in \mathcal{T}, b(t)=\left(\left(a-\delta_{\emptyset}\right) \star \omega\right)(t) .
$$

## Lemma

The coefficients $\omega$ satisfy the following relation for all $m$-uplets, $m \geq 2$, of trees $\left(u_{1}, \ldots, u_{m}\right) \in \mathcal{T}^{m}$ :

$$
\sum_{\substack{I \cup J=\{1, \ldots, m\}, I \cap J=\emptyset}} \omega\left(x_{i \in I} u_{i} \circ \prod_{j \in J} u_{j}\right)=0,
$$

with the conventions $u \circ \emptyset=u$ and $\emptyset \circ u=\emptyset$.

## From 1-cuts to multicuts

Let $a \in \operatorname{Alg}(\mathcal{H}, \mathbb{R})$ and $b \in \operatorname{VF}(\mathcal{H}, \mathbb{R})$. Then one has

$$
\begin{array}{r}
\forall o=\left(t_{1} \ldots t_{m}\right) \in \mathcal{A T}, \sum_{t \in C(o)} b(t)=0 \text { iff } \\
\forall o=\left(t_{1} \ldots t_{m}\right) \in \mathcal{A T}, \sum_{k=1}^{m}(-1)^{k+1} \sum_{t \in C_{k}(o)} a(u)=0 .
\end{array}
$$

Consider the case $m=3$, and, for instance, $o=\left(t_{1} t_{2} t_{3}\right)$. We have to compute

$$
b\left(t_{1} \circ t_{2} \circ t_{3}\right)+b\left(t_{2} \circ t_{3} \circ t_{1}\right)+b\left(t_{3} \circ t_{1} \circ t_{2}\right)
$$

in terms of the $a$ 's. Given $\left(p_{1}, p_{2}, p_{3}\right)$ in $\mathcal{P}\left(t_{1}\right) \times \mathcal{P}\left(t_{2}\right) \times \mathcal{P}\left(t_{3}\right)$, a partition $p \in \mathcal{P}\left(t_{1} \circ t_{2} \circ t_{3}\right)$ is of the form

Table: Terms in the substitution law for $t_{1} \circ t_{2} \circ t_{3}$

| $p$ | $p_{1} \circ p_{2} \circ p_{3}$ | $p_{1} \circ p_{2} \diamond p_{3}$ | $p_{1} \diamond p_{2} \circ p_{3}$ | $p_{1} \diamond p_{2} \diamond p_{3}$ |
| :---: | :---: | :---: | :---: | :---: |
| $\chi_{p}$ | $\chi p_{1} \times \chi p_{2} \times \chi_{p_{3}}$ | $\chi p_{1} \times \chi p_{2} \circ \chi_{p_{3}}$ | $\chi p_{1} \circ \chi p_{2} \times \chi p_{3}$ | $\chi p_{1} \circ \chi_{p_{2}} \circ \chi p_{3}$ |
| $v_{p}$ | $v_{p_{1}}^{*} v_{p_{2}}^{*} v_{p_{3}}^{*} r_{p_{1}} \circ r_{p_{2}} \circ r_{p_{3}}$ | $v_{p_{1}}^{*} v_{p_{2}}^{*} v_{p_{3}}^{*}\left(r_{p_{1}} \circ r_{p_{2}}\right) r_{p_{3}}$ | $v_{p_{1}}^{*} v_{p_{2}}^{*} v_{p_{3}}^{*} r_{p_{1}}\left(r_{p_{2}} \circ r_{p_{3}}\right)$ | $v_{p_{1}}^{*} v_{p_{2}}^{*} v_{p_{3}}^{*} r_{p_{1}} r_{p_{2}} r_{p_{3}}$ |

Hence,

$$
\begin{array}{r}
b\left(t_{1} \circ t_{2} \circ t_{3}\right)=\sum_{\left(p_{1}, p_{2}, p_{3}\right)} a\left(v_{p_{1}}^{*} v_{p_{2}}^{*} v_{p_{3}}^{*}\right) \times \\
\left(\omega\left(\chi_{p_{1}} \times \chi_{p_{2}} \times \chi_{p_{3}}\right) a\left(r_{p_{1}} \circ r_{p_{2}} \circ r_{p_{3}}\right)\right. \\
1 \text {-cut term } \\
+\omega\left(\chi_{p_{1}} \times \chi_{p_{2}} \circ \chi_{p_{3}}\right) a\left(r_{p_{1}} \circ r_{p_{2}}\right) a\left(r_{p_{3}}\right) \\
+\omega\left(\chi_{p_{1}} \circ \chi_{p_{2}} \times \chi_{p_{3}}\right) a\left(r_{p_{1}}\right) a\left(r_{p_{2}} \circ r_{p_{3}}\right) \\
\text { 2-cut term } \\
+\omega\left(\chi_{p_{1}} \circ \chi_{p_{2}} \circ \chi_{p_{3}}\right) a\left(r_{p_{1}}\right) a\left(r_{p_{2}}\right) a\left(r_{p_{3}}\right) \\
\text { 3-cut term })
\end{array}
$$

For $b\left(t_{1} \circ t_{2} \circ t_{3}\right)+b\left(t_{2} \circ t_{3} \circ t_{1}\right)+b\left(t_{3} \circ t_{1} \circ t_{2}\right)$ we get:
1-cut terms

$$
\omega\left(\chi_{p_{1}} \times \chi_{p_{2}} \times \chi_{p_{3}}\right)\left(a\left(r_{p_{1}} \circ r_{p_{2}} \circ r_{p_{3}}\right)+a\left(r_{p_{2}} \circ r_{p_{3}} \circ r_{p_{1}}\right)+a\left(r_{p_{3}} \circ r_{p_{1}} \circ r_{p_{2}}\right)\right) .
$$

2-cut terms

$$
a\left(r_{p_{3}}\right) a\left(r_{p_{1}} \circ r_{p_{2}}\right)\left(\omega\left(\chi_{p_{1}} \times \chi_{p_{2}} \circ \chi_{p_{3}}\right)+\omega\left(\chi_{p_{3}} \circ \chi_{p_{1}} \times \chi_{p_{2}}\right)\right)+\ldots
$$

where

$$
\omega\left(\chi_{p_{1}} \times \chi_{p_{2}} \circ \chi_{p_{3}}\right)+\omega\left(\chi_{p_{3}} \circ \chi_{p_{1}} \times \chi_{p_{2}}\right)=-\omega\left(\chi_{p_{1}} \times \chi_{p_{2}} \times \chi_{p_{3}}\right)
$$

3-cut terms

$$
\begin{aligned}
& a\left(r_{p_{1}}\right) a\left(r_{p_{2}}\right) a\left(r_{p_{3}}\right)\left(\omega\left(\chi_{p_{1}} \circ \chi_{p_{2}} \circ \chi_{p_{3}}\right)+\omega\left(\chi_{p_{2}} \circ \chi_{p_{3}} \circ \chi_{p_{1}}\right)+\omega\left(\chi_{p_{3}} \circ \chi_{p_{1}} \circ \chi_{p_{2}}\right)\right) \\
& \text { i.e., } a\left(r_{p_{1}}\right) a\left(r_{p_{2}}\right) a\left(r_{p_{3}}\right) \omega\left(\chi_{p_{1}} \times \chi_{p_{2}} \times \chi_{p_{3}}\right) .
\end{aligned}
$$

## THIS IS THE END

