

# On quantization of quadratic Poisson structures

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**Abstract :** Any classical r-matrix on the Lie algebra of linear operators on a real vector space  $V$  gives rise to a quadratic Poisson structure on  $V$  which admits a deformation quantization stemming from the construction of V. Drinfel'd [Dr], [Gr]. We exhibit in this article an example of quadratic Poisson structure which does not arise this way.

## I. Introduction

let  $V$  be a finite-dimensional real vector space. The linear action of the Lie group  $Gl(V)$  on  $V$  induces by differentiation a Lie algebra isomorphism from  $\mathfrak{g} = \mathfrak{gl}(V)$  to the Lie algebra of linear vector fields on  $V$ . Given a basis  $(e_1, \dots, e_n)$  and then identifying  $\mathfrak{gl}(V)$  with the Lie algebra of real  $n \times n$  matrices the isomorphism is given by :

$$J(E_{ij}) = x_i \partial_j,$$

where  $E_{ij}$  is the matrix with entries all vanishing except one equal to 1 on the  $i^{\text{th}}$  line and  $j^{\text{th}}$  column.

There is a unique way to extend the Lie bracket of  $\mathfrak{g}$  to a graded Lie bracket, called the *Schouten bracket* on the shifted exterior algebra  $\Lambda(\mathfrak{g})[1]$  in a way compatible with the exterior product. The shift means that elements of  $\Lambda^k(\mathfrak{g})$  are of degree  $k - 1$ , and then the Schouten bracket maps  $\Lambda^k(\mathfrak{g}) \times \Lambda^l(\mathfrak{g})$  to  $\Lambda^{k+l-1}(\mathfrak{g})$ . The exterior algebra  $\Lambda(\mathfrak{g})$  inherits then a structure of Gerstenhaber algebra (cf. for example [V], introduction).

The space  $T^{\text{poly}}(V)$  of polyvector fields on  $M$  is also endowed with a Gerstenhaber algebra structure, with Schouten bracket extending Lie bracket of vector fields [V]. The subalgebra (for exterior product) generated by linear vector fields is a Gerstenhaber subalgebra  $\tilde{\Lambda}(V)$  of  $T^{\text{poly}}(V)$ . The isomorphism  $J$  extends to a surjective Gerstenhaber algebra morphism :

$$J^\bullet : \Lambda^\bullet(\mathfrak{g}) \mapsto \tilde{\Lambda}^\bullet(V).$$

Map  $J^k$  has nontrivial kernel for  $k \geq 2$  as long as  $V$  has dimension  $\geq 2$  : for example we have :

$$J^2(E_{ij} \wedge E_{kj}) = 0.$$

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A *classical r-matrix* on  $\mathfrak{g}$  is by definition an element  $r$  of  $\mathfrak{g} \wedge \mathfrak{g}$  such that  $[r, r] = 0$ . According to the discussion above the bivector field  $J^2(r)$  defines then a quadratic Poisson structure on  $V$ . A natural question arises then : can one recover this way any quadratic Poisson structure on  $V$ ? It was claimed true in [BR] but Z.H. Liu and P. Xu discovered that the authors' argument was not correct [LX]. They brought up a positive answer in the two-dimensional case ([LX] prop. 2.1) : namely the general quadratic Poisson structure  $(ax_1^2 + 2bx_1x_2 + cx_2^2)\partial_1 \wedge \partial_2$  is equal to  $J^2(r)$  with :

$$r = \begin{pmatrix} b & -a \\ c & -b \end{pmatrix} \wedge \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}.$$

We give here a *negative* answer to this question in general : after a somewhat lengthy but elementary computation we show in § III that bivector field  $(x_1^2 + x_2x_3)\partial_2 \wedge \partial_3$  on  $\mathbb{R}^3$  is a counterexample to this conjecture : it is outside the image of the set of r-matrices by  $J^2$ .

We recall in § II the construction by V.G. Drinfel'd of a translation-invariant deformation quantization on any Lie group  $G$  once given a classical r-matrix on the Lie algebra  $\mathfrak{g}$  [Dr], [T]. The problem reduces to the case when  $r$  is non-degenerate, and the deformation quantization is then obtained by suitable restriction and transportation of Baker-Campbell-Hausdorff deformation quantization ([Ka]) of the dual  $\tilde{\mathfrak{g}}^*$  of the central extension  $\tilde{\mathfrak{g}}$  of  $\mathfrak{g}$  defined by  $r$ . The construction works moreover for any Kontsevich-type star product [ABM] on  $\tilde{\mathfrak{g}}^*$ . For  $\mathfrak{g} = \mathfrak{gl}(V)$ , such a star product on  $\tilde{\mathfrak{g}}^*$  gives almost immediately through this construction a deformation quantization of quadratic Poisson structure  $J^2(r)$  on  $V$ .

Let us mention that the existence of a deformation quantization for any quadratic Poisson structure has been proved by Omori, Maeda and Yoshioka [OMY], a few years before the proof by Kontsevich of the existence of a deformation quantization for any Poisson structure on any manifold [K], [CFT]. Explicit computations for all quadratic Poisson structures in dimension 3 (then including our counterexample as well) have been performed by El Galiou and Tihami [ET], by a case-by-case method based on the classification of Dufour and Haraki [DH].

## II. Quantization of Poisson structures coming from r-matrices

Let  $\mathfrak{g}$  be a Lie algebra, and let  $r \in \mathfrak{g} \wedge \mathfrak{g}$  a classical r-matrix. It defines an antisymmetric operator :

$$\tilde{r} : \mathfrak{g}^* \longrightarrow \mathfrak{g}.$$

Classical Yang-Baxter equation  $[r, r] = 0$  is equivalent to :

$$\langle \xi, [\tilde{r}(\eta), \tilde{r}(\zeta)] \rangle + \langle \eta, [\tilde{r}(\zeta), \tilde{r}(\xi)] \rangle + \langle \zeta, [\tilde{r}(\xi), \tilde{r}(\eta)] \rangle = 0 \quad (*)$$

for any  $\xi, \eta, \zeta \in \mathfrak{g}^*$ . The r-matrix  $r$  defines a left translation-invariant Poisson structure on any Lie group  $G$  with Lie algebra  $\mathfrak{g}$ .

### II.1. A central extension

We can firstly suppose  $r$  nondegenerate, i.e. that  $\tilde{r}$  is inversible with inverse  $\tilde{\omega}$ , where  $\omega$  belongs to  $\mathfrak{g}^* \wedge \mathfrak{g}^*$ . Classical Yang-Baxter equation is in this case equivalent to :

$$\langle \tilde{\omega}X, [Y, Z] \rangle + \langle \tilde{\omega}Y, [Z, X] \rangle + \langle \tilde{\omega}Z, [X, Y] \rangle = 0$$

for any  $X, Y, Z \in \mathfrak{g}$ , i.e is equivalent to the fact that  $\omega$  is a 2-cocycle with values in the trivial representation. Let  $\tilde{\mathfrak{g}}$  the central extension of  $\mathfrak{g}$  by this cocycle, defined by  $\tilde{\mathfrak{g}} = \mathfrak{g} \oplus \mathbb{R}$  with bracket :

$$[X + \alpha, Y + \beta] = [X, Y] + \langle \omega, X \wedge Y \rangle.$$

The cocycle condition on  $\omega$  is equivalent to de Jacobi identity for this bracket. Let  $X_0 = (0, 1) \in \tilde{\mathfrak{g}}$ , and let  $\mathcal{H}$  the hyperplane in  $\tilde{\mathfrak{g}}^*$  defined by :

$$\mathcal{H} = \{ \xi \in \tilde{\mathfrak{g}}^* / \langle \xi, X_0 \rangle = 1 \}.$$

It is the symplectic leaf through the point  $\xi_0$  defined by  $\langle \xi_0, \mathfrak{g} \rangle = 0$  and  $\langle \xi_0, X_0 \rangle = 1$ . It is then the coadjoint orbit  $\text{Ad}^* \tilde{G} \cdot \xi_0$  for any Lie group  $\tilde{G}$  with Lie algebra  $\tilde{\mathfrak{g}}$ .

## II.2. Kontsevich star products (after [ABM])

The linear Poisson manifold  $\tilde{\mathfrak{g}}^*$  admits a whole bunch of equivalent  $\text{Ad}^* \tilde{G}$ -invariant deformation quantizations which can be built from the enveloping algebra  $\mathcal{U}(\tilde{\mathfrak{g}})$ , for example the Baker-Campbell-Hausdorff quantization or the Kontsevich quantization [K], [ABM], [Ka], [Di]. The baker-Campbell-Hausdorff quantization is given by the following integral formula [ABM] :

$$(u \#_{BCH} v)(\xi) = \iint_{\mathfrak{g} \times \mathfrak{g}} \mathcal{F}^{-1}u(x)\mathcal{F}^{-1}v(y) e^{i\langle \xi, x \cdot y \rangle_{\hbar}} dx dy,$$

where the inverse Fourier transform is given by :

$$\mathcal{F}^{-1}u(x) = (2\pi)^{-n} \int_{\mathfrak{g}^*} u(\eta) e^{i\langle x, \eta \rangle} d\eta,$$

and  $x \cdot y$  stands for the Baker-Campbell-Hausdorff expansion :

$$x + y + \frac{\hbar}{2}[x, y] + \frac{\hbar^2}{12}([x, [x, y]] + [y, [y, x]]) + \dots$$

The Lebesgue measure  $d\eta$  on  $\mathfrak{g}^*$  is normalized so that it is the dual measure of Lebesgue measure  $dx$  on  $\mathfrak{g}$ . The quantizations we can consider here are the ones called ‘‘Kontsevich star products’’ in [ABM]. They are all equivalent to the BCH quantization. The equivalence is a formal series of differential operators with constant coefficients on  $\mathfrak{g}^*$  precisely given by a formal series of  $G$ -invariant polynomials on  $\mathfrak{g}$  of the following form :

$$F(x) = 1 + \sum_{k \geq 1} \hbar^{2k} \sum_{c \geq 1} \sum_{(s_1, \dots, s_c) \in S_{2k}^c} a_{s_1, \dots, s_c} \text{Tr}(\text{ad } x)^{s_1} \dots \text{Tr}(\text{ad } x)^{s_c},$$

where  $S_{2k}^c$  stands for those  $(s_1, \dots, s_c)$  in  $\mathbb{N}^c$  such that  $s_1 + \dots + s_c = 2k$ ,  $s_1 \leq s_2 \leq \dots \leq s_c$  and  $s_j \neq 1$ . The star product obtained this way admits the following integral form :

$$(u \# v)(\xi) = \iint_{\mathfrak{g} \times \mathfrak{g}} \mathcal{F}^{-1}u(x) \mathcal{F}^{-1}v(y) \frac{F(-ix)F(-iy)}{F(-i(x \cdot y))} e^{i \langle \xi, x \cdot y \rangle} dx dy.$$

### II.3. Quantization of left-invariant Poisson structures

It is easy to derive from the fact that  $X_0$  is central that any of the deformation quantizations defined above does define by restriction a deformation quantization of  $\mathcal{H}$ . Let  $G$  be the subgroup of  $\tilde{G}$  with Lie algebra  $\mathfrak{g}$ . We clearly have :

$$\text{Ad}^* G \cdot \xi_0 = \text{Ad}^* \tilde{G} \cdot \xi_0 = \mathcal{H}.$$

It is moreover easy to check that the stabilizer of  $\xi_0$  in  $\tilde{G}$  is the one-dimensional subgroup with Lie algebra generated by  $X_0$ . It is a simple consequence of the nondegeneracy of the alternate bilinear form  $\omega$ . The dimension of  $G$  is the equal to the dimension of  $\mathcal{H}$ . The map :

$$\begin{aligned} \varphi : G' &\longrightarrow \mathcal{H} \\ g &\longmapsto \text{Ad}^* g \cdot \xi_0 \end{aligned}$$

is then a local  $G$ -equivariant diffeomorphism near the identity (with left translation on the left-hand side and coadjoint action on the right-hand side). We can then transport any deformation quantization of  $\mathcal{H}$  and get a left translation-invariant deformation quantization of a neighbourhood of the identity in  $G$ . It extends by translation invariance to the whole group  $G$ , as well as to any Lie group  $G'$  locally isomorphic to  $G$ .

The deformation quantization on  $G$  can be written :

$$u \# v = \sum_{k \geq 0} \hbar^k C_k(u, v),$$

where the  $C_k$ 's are left-invariant bidifferential operators on  $G$ . There exists then an element  $F = \sum \hbar^k F_k$  in  $(\mathcal{U}(\mathfrak{g}) \otimes \mathcal{U}(\mathfrak{g}))[[\hbar]]$  such that :

$$u \# v(g) = F(u \otimes v)(g, g). \quad (**)$$

Let us now fix a basis  $x_1, \dots, x_n$  of  $\mathfrak{g}$ , and consider elements of  $\mathcal{U}(\mathfrak{g})$  as polynomials  $F(x) = F(x_1, \dots, x_n)$  of the  $n$  noncommuting variables  $x_1, \dots, x_n$ , which satisfy the relations :

$$x_i x_j - x_j x_i = [x_i, x_j] = \sum_k c_{ij}^k x_k.$$

Introducing a second identical set of noncommuting variables  $y = (y_1, \dots, y_n)$  commuting with the  $x_j$ 's we can write any element  $A \in \mathcal{U}(\mathfrak{g}) \otimes \mathcal{U}(\mathfrak{g})$  as  $A(x, y)$ . The element  $F$  defined above can then be written  $F(x, y)$  as a formal series with coefficients  $F_k(x, y)$ .

**Proposition II.1.**

The formal series  $F = F(x, y) \in (\mathcal{U}(\mathfrak{g}) \otimes \mathcal{U}(\mathfrak{g}))[[\hbar]]$  above verifies :

- 1)  $F_0(x, y) = 0$ .
- 2)  $F_1(x, y) = \frac{1}{2} \sum_{i,j} r_{ij} x_i y_j$ .
- 3)  $F_k(x, 0) = F_k(0, y) = 0$  for  $k \geq 1$ .
- 4)  $F(x + y, z)F(x, y) = F(x, y + z)F(y, z)$ .

Conversely any  $F(x, y)$  endowed with those 4 properties defines by formula (\*\*) a left translation deformation quantization of  $G$ .

*Proof.* It is well-known : see for example [Dr], [T] : first condition comes from the fact that  $C_0(u, v)$  is the ordinary product  $uv$ . Second property comes from the expression of left-invariant Poisson bracket on  $G$  defined from the r-matrix, third property expresses the fact that  $1\#u = u\#1 = u$ , and last property is an expression of the associativity of star product  $\#$ . Let us elaborate a bit on that last point : any element  $X_j$  of the basis corresponds to polynomial expression  $G(x) = x_j$ . The Leibniz rule :

$$X_j.(\varphi\psi) = (X_j.\varphi)\psi + \varphi.(X_j.\psi)$$

can be written as :

$$G(x) \circ m = m \circ G(x + y),$$

where  $m : C^\infty(G \times G) \rightarrow C^\infty(G)$  stands for multiplication, here the restriction to the diagonal. The formula above extends to any polynomial expression  $G$  representing any element of the enveloping algebra. We have then :

$$\begin{aligned} (u\#v)\#w &= (m \circ F(u \otimes v))\#w \\ &= m \circ F\left((m \circ F(u \otimes v)) \otimes w\right) \\ &= m \circ F \circ (m \otimes I) \circ (F \otimes I)(u \otimes v \otimes w) \\ &= m \circ F(x, z) \circ (m \otimes I) \circ F(x, y)(u \otimes v \otimes w) \\ &= m \circ (m \otimes I) \circ F(x + y, z)F(x, y)(u \otimes v \otimes w). \end{aligned}$$

Similarly we have :

$$u\#(v\#w) = m \circ (I \otimes m) \circ F(x, y + z)F(y, z)(u \otimes v \otimes w).$$

The associativity condition for product  $\#$  is then equivalent to property 4) of the proposition. •

let us now look at the case when  $r$  is degenerate. Then the image  $\mathfrak{g}_0$  of  $\tilde{r}$  is a subspace strictly contained in  $\mathfrak{g}$ . By skew-symmetry  $\mathfrak{g}_0$  is also the orthogonal of the kernel of  $\tilde{r}$ , and classical Yang-Baxter equation  $[r, r] = 0$  ensures thanks to (\*) that  $\mathfrak{g}_0$  is a Lie subalgebra of  $\mathfrak{g}$ . We get this way a nondegenerate  $r_0 \in \mathfrak{g}_0 \wedge \mathfrak{g}_0$  such that  $[r_0, r_0] = 0$ . Applying the procedure above we get an  $F = \sum \hbar^k F_k$  in  $(\mathcal{U}(\mathfrak{g}_0) \otimes \mathcal{U}(\mathfrak{g}_0))[[\hbar]]$  which can be seen as an element of  $(\mathcal{U}(\mathfrak{g}) \otimes \mathcal{U}(\mathfrak{g}))[[\hbar]]$ .

#### II.4. A class of easily quantizable Poisson structures

Let  $G$  be a Lie group with Lie algebra  $\mathfrak{g}$ . Let  $F(x, y)$  a formal series in  $\mathcal{U}(\mathfrak{g}) \otimes \mathcal{U}(\mathfrak{g})[[\hbar]]$  satisfying properties 1-4 of proposition II.1 (for example that one constructed from an r-matrix along the lines above). Let  $M$  be any differentiable manifold endowed with an action of  $G$ . The differentiation of this action induces a Lie algebra morphism from  $\mathfrak{g}$  to the vector fields on  $M$ , which extends to an algebra morphism from  $\mathcal{U}(\mathfrak{g})$  to the algebra of differential operators on  $M$ . Similarly it induces an algebra morphism from  $\mathcal{U}(\mathfrak{g}) \otimes \mathcal{U}(\mathfrak{g})$  to the algebra of differential operators on  $M \times M$ . The formal series of bidifferential operators defined by the formula :

$$* = m \circ F(x, y)$$

(where  $m : C^\infty(M \times M) \rightarrow C^\infty(M)$  stands for ordinary multiplication of functions on  $M$ ) defines then a star product on  $M$ , the associated Poisson bivector being defined by  $F_1(x, y) - F_1(y, x)$ . The proof of this fact goes the same way as that of proposition II.1. It is easily seen that if  $F(x, y)$  comes from a classical r-matrix  $r \in \mathfrak{g} \wedge \mathfrak{g}$  then the Poisson structure on  $M$  is  $J^2(r)$  where  $J^\bullet$  is the Gerstenhaber algebra morphism from  $\Lambda(\mathfrak{g})$  to multivector fields on  $M$  extending the action of  $\mathfrak{g}$ .

We will be interested in the sequel by the following particular situation : the manifold  $M$  is a vector space  $V$ , the action of  $G$  is linear, and there is a classical r-matrix  $r$  on  $\mathfrak{g}$ . We can as in the introduction view  $J$  as a Lie algebra morphism from  $\mathfrak{g}$  to the space of linear vector fields on  $V$ , and extend  $J$  to a morphism  $J^\bullet$  of Gerstenhaber algebras from  $\Lambda(\mathfrak{g})$  to  $\tilde{\Lambda}(V)$ . In particular  $J^2(r)$  defines a quadratic Poisson structure on  $V$ , and formula just above gives a quantization of this particular quadratic Poisson structure.

### III. Quadratic Poisson structures and r-matrices

#### III.1. Some definitions

We keep the notations of the introduction. The Gerstenhaber algebra  $\tilde{\Lambda}(V)$  can be written as :

$$\tilde{\Lambda}(V) = \bigoplus_{n \geq 0} (S^n(V) \otimes \Lambda^n(V))[1] = \bigoplus_{n \geq 0} \tilde{\Lambda}^n(V)[1].$$

A quadratic Poisson structure on  $V$  can be defined as a bivector field  $\Lambda$  in  $\tilde{\Lambda}^2(V)$  such that :

$$[\Lambda, \Lambda] = 0.$$

Let  $\Lambda$  be an element of  $\tilde{\Lambda}^2(V)$ , and let  $r$  an element of  $\Lambda^2(\mathfrak{g})$  such that  $J^2(r) = \Lambda$ . It is then obvious that  $[\Lambda, \Lambda] = 0$  if and only if  $J^3([r, r]) = 0$ . If  $n \geq 2$  then  $J^2$  and  $J^3$  have

nontrivial kernels : Precisely we have :

$$\dim \ker J^2 = \frac{n^2(n^2 - 1)}{4} \quad \text{and} \quad \dim \ker J^3 = \frac{n^2(n^2 - 1)(5n^2 - 8)}{36}.$$

### III.2. A counterexample in dimension 3

With the notations of § I, an element of  $\mathfrak{g} \wedge \mathfrak{g}$  can be written as :

$$r = \sum_{i,j,k,l=1}^n r_{ik}^{jl} E_{ij} \wedge E_{kl}.$$

We shall need for further calculations the following result :

#### Proposition III.1.

Let  $r = \sum_{i,j,k,l=1}^n r_{ik}^{jl} E_{ij} \wedge E_{kl}$  be an element of  $\mathfrak{g} \wedge \mathfrak{g}$  then  $[r, r] = 0$  if and only if for any  $i, j, k, l, m, p \in \{1, \dots, n\}$  such that  $(i, j) < (k, l) < (m, p)$  according to lexicographical order we have :

$$\sum_{d=1}^n r_{ik}^{dl} r_{dm}^{jp} - r_{mk}^{dl} r_{di}^{pj} + r_{km}^{dp} r_{di}^{lj} - r_{im}^{dp} r_{dk}^{jl} + r_{mi}^{dj} r_{dk}^{pl} - r_{ki}^{dj} r_{dm}^{lp} = 0.$$

*Proof.* This proposition is a direct consequence of formula :

$$[E_{ij}, E_{kl}] = \delta_{jk} E_{il} - \delta_{li} E_{kj}$$

and the following lemma :

#### Lemma III.1.

Let  $\mathfrak{h}$  be a finite-dimensional Lie algebra and let  $X_1, \dots, X_N$  be a basis of  $\mathfrak{h}$ .

If  $r = \sum_{I,J=1}^N r^{IJ} X_I \wedge X_J$  is an element of  $\mathfrak{h} \wedge \mathfrak{h}$  ( $r^{IJ} = -r^{JI}$ ) then

$$[r, r] = 4 \sum_{I,J,K,L=1}^N r^{IJ} r^{KL} [X_I, X_K] \wedge X_J \wedge X_L.$$

*Remark :* We can directly show proposition III.1 using relation (\*) of beginning of § II applied to elements of the dual basis of  $X_1, \dots, X_n$ . •

**Proposition III.2.**

The Poisson structure  $\Lambda$  on  $\mathbb{R}^3$  given by

$$\Lambda = (x_1^2 + \alpha x_2 x_3) \partial_2 \wedge \partial_3$$

with  $\alpha \neq 0$  is not the image of a classical  $r$ -matrix by  $J^2$ .

*Proof.* An element  $r = \sum_{i,j,k,l=1}^n r_{ik}^{jl} E_{ij} \wedge E_{kl}$  of  $\mathfrak{g} \wedge \mathfrak{g}$  is parametrized by 36 coefficients  $r_{ik}^{jl}$ .

It has image  $\Lambda$  if and only if the following 18 equations are satisfied :

$$\begin{aligned} r_{11}^{12} = r_{11}^{13} = r_{22}^{12} = r_{22}^{13} = r_{22}^{23} = r_{33}^{12} = r_{33}^{13} = r_{33}^{23} = 0, & \quad r_{11}^{23} = 1 \\ r_{12}^{12} = r_{12}^{21}, & \quad r_{12}^{13} = r_{12}^{31}, & \quad r_{13}^{12} = r_{13}^{21} \\ r_{13}^{13} = r_{13}^{31}, & \quad r_{12}^{23} = r_{12}^{32}, & \quad r_{13}^{23} = r_{13}^{32} \\ r_{23}^{12} = r_{23}^{21}, & \quad r_{23}^{13} = r_{23}^{31}, & \quad r_{23}^{32} = r_{23}^{23} - \alpha \end{aligned}$$

To lighten writing we rename the 18 remaining unknowns as follows :

$$\begin{aligned} r_{12}^{11} = a, & \quad r_{12}^{12} = b, & \quad r_{12}^{13} = c, & \quad r_{13}^{11} = d, & \quad r_{13}^{12} = e, & \quad r_{13}^{13} = f, \\ r_{12}^{22} = g, & \quad r_{12}^{23} = h, & \quad r_{13}^{22} = i, & \quad r_{13}^{23} = j, & \quad r_{12}^{33} = k, & \quad r_{13}^{33} = l, \\ r_{23}^{11} = m, & \quad r_{23}^{12} = n, & \quad r_{23}^{13} = p, & \quad r_{23}^{22} = q, & \quad r_{23}^{23} = r, & \quad r_{23}^{33} = s. \end{aligned}$$

The unknown  $r$  must not be confused with the classical  $r$ -matrix on the whole. The context will not lead to any confusion. We have then :

$$\begin{aligned} r_{11}^{23} &= 1, \\ r_{12}^{21} = b, & \quad r_{13}^{21} = e, & \quad r_{12}^{31} = c, \\ r_{12}^{32} = h, & \quad r_{13}^{31} = f, & \quad r_{13}^{32} = j, \\ r_{23}^{21} = n, & \quad r_{23}^{31} = p, & \quad r_{23}^{32} = r - \alpha, \end{aligned}$$

and other  $r_{ik}^{jl}$  are equal to 0.

If  $r$  is such an element the equation  $[r, r] = 0$  develops according to Proposition III.2. into a system of 84 equations involving our 18 unknowns  $a, b, \dots, s$ , given by the vanishing of the 84 coefficients of elements of the basis  $E_{ij} \wedge E_{kl} \wedge E_{mn}$  of  $\mathfrak{g} \wedge \mathfrak{g} \wedge \mathfrak{g}$ . The 84 equations reduce to 66 thanks to the fact that we already have  $[\Lambda, \Lambda] = 0$ . But we shall only consider 20 of them, which will be sufficient for exhibiting the counterexample :

Let us order the  $E_{ij}$ 's lexicographically from first to 9th, rename them accordingly ( $A_1 = E_{11}, A_2 = E_{12}, \dots, A_9 = E_{33}$ ), and labelize by  $(x, y, z)$  the equation obtained by the vanishing of the coefficient of  $A_x \wedge A_y \wedge A_z$ . We shall consider precisely the following equations :



$$\begin{array}{ll}
(1, 2, 5) \quad ci - eh + n = 0 & (1, 2, 9) \quad eh - ci + d = 0 \\
(1, 3, 5) \quad cj - ek - a = 0 & (1, 3, 7) \quad ce + f^2 - ja - ld + m = 0 \\
(1, 3, 9) \quad ek - cj + p = 0 & (1, 4, 5) \quad mh - nc = 0 \\
(1, 4, 6) \quad mk - pc = 0 & (1, 5, 6) \quad \alpha c + nk - ph = 0 \\
(1, 7, 8) \quad en - im = 0 & (1, 7, 9) \quad ep - jm = 0 \\
(1, 8, 9) \quad \alpha e + ip - jn = 0 & (2, 3, 9) \quad -3f + r + ik - jh = 0 \\
(2, 4, 5) \quad ag - b^2 + nh - qc = 0 & (2, 5, 6) \quad 2\alpha h + bh - rh + qk - cg = 0 \\
(2, 8, 9) \quad ej + ir + \alpha i - jq - fi = 0 & (3, 8, 9) \quad is + 2\alpha j + el - fj - jr = 0 \\
(4, 5, 7) \quad pn - rm - bm + na = 0 & (5, 6, 9) \quad -nk - rs + hp + s(r - \alpha) = 0 \\
(5, 8, 9) \quad nj - ip + \alpha q = 0 & (6, 8, 9) \quad ln - pj + qs - (r - \alpha)^2 = 0.
\end{array}$$

Consider the following two sums :

$$\begin{aligned}
(1, 2, 5) + (1, 2, 9) : \quad n + d &= 0 \\
(1, 3, 5) + (1, 3, 9) : \quad p - a &= 0.
\end{aligned}$$

Hence  $n = -d$  and  $p = a$ . We will discuss the four cases  $a = d = 0$ ,  $a = 0$  and  $d \neq 0$ ,  $a \neq 0$  and  $d = 0$ ,  $a \neq 0$  and  $d \neq 0$ .

**First case** :  $a = d = 0$ . Then looking successively at the following equations we get :

$$\begin{aligned}
(5, 8, 9) \implies q &= 0 & (6, 8, 9) \implies r &= \alpha & (1, 5, 6) \implies c &= 0 \\
(1, 8, 9) \implies e &= 0 & (2, 4, 5) \implies b &= 0 & (2, 5, 6) \implies h &= 0 \\
(4, 5, 7) \implies m &= 0 & (5, 6, 9) \implies s &= 0 & (1, 3, 7) \implies f &= 0 \\
(3, 8, 9) \implies j &= 0 & (2, 8, 9) \implies i &= 0 & (2, 3, 9) \implies \alpha &= 0,
\end{aligned}$$

hence a contradiction to the hypothesis  $\alpha \neq 0$ .

**Second case** :  $a = 0$  and  $d \neq 0$  (hence  $n \neq 0$ ).

$$(1, 4, 6) \implies mk = 0.$$

*First subcase* :  $m = 0$ . Then :

$$(1, 4, 5) \implies c = 0 \quad (1, 7, 8) \implies e = 0 \quad (1, 2, 5) \implies n = 0,$$

hence a contradiction.

*Second subcase* :  $m \neq 0$ , hence  $k = 0$ .

$$(1, 5, 6) \implies c = 0 \quad (1, 4, 5) \implies h = 0 \quad (1, 2, 5) \implies n = 0,$$

hence a contradiction again.

**Third case** :  $a \neq 0$  and  $d = 0$  (hence  $p \neq 0$ ).

$$(1, 4, 5) \implies mh = 0.$$

First subcase :  $m = 0$ . Then :

$$(1, 4, 6) \implies c = 0 \quad (1, 7, 9) \implies e = 0 \quad (1, 3, 5) \implies a = 0,$$

hence a contradiction.

Second subcase :  $m \neq 0$ , hence  $h = 0$ .

$$(1, 5, 6) \implies c = 0 \quad (1, 4, 6) \implies k = 0 \quad (1, 3, 5) \implies a = 0,$$

hence a contradiction again.

**Fourth case** :  $a \neq 0$  and  $d \neq 0$ .

First subcase :  $m = 0$ .

$$(1, 4, 5) \implies c = 0 \quad (1, 7, 8) \implies e = 0 \quad (1, 3, 5) \implies a = 0,$$

contradiction.

Second subcase :  $m \neq 0$ .

$$(1, 4, 6) \implies k = \frac{a}{m}c \quad (1, 7, 9) \implies j = \frac{a}{m}e \quad (1, 3, 5) \implies a = 0,$$

contradiction.

This proves proposition III.2. •

### III.3. Cartan-type quadratic Poisson structures

Recall from [DH] that the curl of a Poisson structure  $\Lambda = \sum_{i,j} \Lambda^{ij} \partial_i \wedge \partial_j$  is defined by :

$$\text{rot } \Lambda = \sum_{i,j} \partial_j \Lambda^{ij} \cdot \partial_i.$$

It is a linear vector field (and hence can be viewed as an  $n \times n$  matrix) when  $\Lambda$  is quadratic. A quadratic Poisson structure is of *Cartan type* if it can be written for some choice of coordinates as :

$$\Lambda = \sum_{i,j=1}^n c_{ij} x_i x_j \partial_i \wedge \partial_j$$

with  $c_{ji} = -c_{ij}$ . J.P. Dufour and A. Haraki proved the following result :

**Theorem III.3** (Dufour - Haraki).

Any quadratic Poisson structure the curl of which has eigenvalues  $\lambda_i$  such that  $\lambda_i + \lambda_j \neq \lambda_r + \lambda_s$  for any  $(i, j, r, s)$  with  $r \neq s$  and  $\{i, j\} \neq \{r, s\}$  is of Cartan type.

Such a Cartan-type Poisson structure is image by  $J^{(2)}$ , of a classical  $r$ -matrix, namely :

$$r = \sum_{i,j} c_{ij} E_{ii} \wedge E_{jj}.$$

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