A stabilized finite element projection scheme for incompressible fluid flow

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Stokes Equations

Consider the time-dependent Stokes equations :

$\frac{\partial \mathbf{u}}{\partial t} - \Delta \mathbf{u} + \nabla p = \mathbf{f}$	in $\Omega \times (0, T)$
$\operatorname{div} \mathbf{u} = 0$	in $\Omega imes (0, T)$
u = 0	on $\Gamma \times (0, T)$

where $\Omega \subset \mathbb{R}^d$, $\Gamma = \partial \Omega$, $\mathbf{f} \in L^2(\Omega)^d$.

A projection scheme

$$\begin{cases} \frac{\widetilde{\mathbf{u}}^{n+1} - \mathbf{u}^n}{\delta t} - \Delta \widetilde{\mathbf{u}}^{n+1} = \mathbf{f}^{n+1} - \nabla p^n & \text{in } \Omega\\ \widetilde{\mathbf{u}}^{n+1} = 0 & \text{on } \Gamma \end{cases}$$

$$\begin{cases} \frac{\mathbf{u}^{n+1} - \widetilde{\mathbf{u}}^{n+1}}{\delta t} + \nabla (p^{n+1} - p^n) = 0 & \text{in } \Omega\\ \operatorname{div} \mathbf{u}^{n+1} = 0 & \text{in } \Omega \end{cases}$$

$$\mathbf{u}^{n+1} \cdot \mathbf{n} = 0$$
 on Γ

Remarks :

- This scheme is of first order $(O(\delta t))$
- We have :

$$\begin{aligned} &-\Delta(p^{n+1}-p^n)=-\frac{1}{\delta t}\operatorname{div}\widetilde{\mathbf{u}}^{n+1} & \text{ in }\Omega\\ &\frac{\partial(p^{n+1}-p^n)}{\partial n}=0 & \text{ on }\Gamma \end{aligned}$$

Then

$$\mathbf{u}^{n+1} = \widetilde{\mathbf{u}}^{n+1} + \delta t \nabla (p^{n+1} - p^n)$$
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$$\frac{1}{\delta t} \mathbf{M}(\widetilde{\mathbf{u}}_h^{n+1} - \mathbf{u}_h^n) + \mathbf{A}\widetilde{\mathbf{u}}_h^{n+1} = \mathbf{b}^{n+1} - \mathbf{B}^T p_h^n$$
$$\frac{1}{\delta t} \mathbf{M}(\mathbf{u}_h^{n+1} - \widetilde{\mathbf{u}}_h^{n+1}) + \mathbf{B}^T (p_h^{n+1} - p_h^n) = 0$$
$$\mathbf{B}\mathbf{u}_h^{n+1} = 0.$$

Whence

$$\mathbf{B}\mathbf{M}^{-1}\mathbf{B}^{\mathsf{T}}(p_h^{n+1}-p_h^n)=\frac{1}{\delta t}\mathbf{B}\widetilde{\mathbf{u}}_h^{n+1}$$

Remarks.

- BM⁻¹B^T is an analog of − div_h ∇_h (not −Δ_h !!). This matrix is sparse if a mass lumping is used. In general, it is less sparse than −Δ_h.
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Stabilization of the projection equation A stabilization of the Darcy equation Implementation A stabilization of the Stokes equations

Stabilized Methods

Aim : Avoid the inf-sup condition by modifying in a consistent way the variationsl formulation

 \Rightarrow Possibility of using arbitrary combinations of velocity and pressure spaces.

 $\mathsf{Example}$: Method of $\operatorname{Hughes}, \operatorname{Balestra}, \operatorname{Franca}$ for steady state Stokes equations :

 $\begin{aligned} (\nabla \mathbf{u}_h, \nabla \mathbf{v}) - (\boldsymbol{p}, \operatorname{div} \mathbf{v}) &= (\mathbf{f}, \mathbf{v}) & \mathbf{v} \in V_h \\ (\operatorname{div} \mathbf{u}_h, q) + \alpha \sum_{\mathcal{T}} h_{\mathcal{T}}^2 ((\nabla \boldsymbol{p}_h, \nabla q)_{\mathcal{T}} \\ &- (\Delta \mathbf{u}_h, \nabla q)_{\mathcal{T}}) = \alpha \sum_{\mathcal{T}} h_{\mathcal{T}}^2 (\mathbf{f}, \nabla q)_{\mathcal{T}} & q \in Q_h \end{aligned}$

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A Stabilisation of the projection equation (Hughes, Masud)

The projection equation is a Darcy equation (porous media) : Projection of $H_0^1(\Omega)^d$ on the space $\{\mathbf{v} \in L^2(\Omega)^d; \text{ div } \mathbf{v} = 0\}$:

Given $\mathbf{u} \in H^1_0(\Omega)^d$, Find $\mathbf{v} \in H(\operatorname{div}; \Omega)$, $\rho \in L^2_0(\Omega)$ such that :

$$\begin{cases} \mathbf{v} + \nabla \mathbf{p} = \mathbf{u} & \text{in } \Omega \\ \text{div } \mathbf{v} = 0 & \text{in } \Omega \\ \mathbf{v} \cdot \mathbf{n} = 0 & \text{on } \Gamma \end{cases}$$

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Mixed Formulation :

 T_h : A mesh of Ω (triangles or tetrahedra).

$$V_h = \{ \mathbf{w} \in H(\operatorname{div}, \Omega); \ \mathbf{w}_{|T} \in RT_0(T), \ T \in \mathcal{T}_h \}$$
$$Q_h = \{ q \in L^2(\Omega); \ q_{|T} \in P_0, \ T \in \mathcal{T}_h, \ \int_{\Omega} q = 0 \}$$
$$RT_0(T) = \{ \mathbf{w} : T \to \mathbb{R}^2; \ \mathbf{w}(\mathbf{x}) = \mathbf{a} + b\mathbf{x}, \ \mathbf{a} \in \mathbb{R}^2, \ b \in \mathbb{R} \} \quad T \in \mathcal{T}_h$$

$$\begin{cases} \mathsf{Find} \ (\mathbf{v}_h, p_h) \in V_h \times Q_h \text{ such that }: \\ (\mathbf{v}_h, \mathbf{w}) - (p_h, \operatorname{div} \mathbf{w}) = (\mathbf{u}, \mathbf{w}) & \forall \ \mathbf{w} \in V_h \\ (\operatorname{div} \mathbf{v}_h, q) = 0 & \forall \ q \in Q_h \end{cases}$$

Remark. An efficient method consists in using a mixed hybrid method (DUBOIS, TOUZANI, ZIMMERMAN). It enables decoupling **v** and *p*

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A stabilization of the Darcy equation. Let $V^{S} = L^{2}(\Omega)^{d}$, $Q^{S} = \{q \in H^{1}(\Omega); \int_{\Omega} q = 0\}$. The stabilized formulation reads :

 $\begin{array}{ll} \mathsf{Find} \ (\mathbf{v}, p) \in V^S \times Q^S \ \mathsf{such that}: \\ (\mathbf{v}, \mathbf{w}) + (\nabla p, \mathbf{w}) = (\mathbf{u}, \mathbf{w}) & \forall \ \mathbf{w} \in V^S, \\ - (\mathbf{v}, \nabla q) + (\nabla p, \nabla q) = (\mathbf{u}, \nabla q) & \forall \ q \in Q^S. \end{array}$

Note that this implies :

 $\mathbf{v} + \nabla p = \mathbf{u} \quad \Rightarrow \quad \operatorname{div} \mathbf{v} + \Delta p = \operatorname{div} \mathbf{u}$ $\operatorname{div} \mathbf{v} - \Delta p = -\operatorname{div} \mathbf{u}$

Whence div $\mathbf{v} = 0$. We also deduce $\mathbf{v} \cdot \mathbf{n} = 0$ on Γ .

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We define the forms

 $\begin{aligned} \mathscr{B}((\mathbf{v}, p); (\mathbf{w}, q)) &= (\mathbf{v}, \mathbf{w}) + (\nabla p, \mathbf{w}) - (\mathbf{v}, \nabla q) + (\nabla p, \nabla q) \\ \mathscr{L}((\mathbf{w}, q)) &= (\mathbf{u}, \mathbf{w}) + (\mathbf{u}, \nabla q) \end{aligned}$

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Discretization.

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Convergence Analysis. We have (MASUD-HUGHES)

 $\| \mathbf{v} - \mathbf{v}_h \|_0 + \|
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 $\|\mathbf{v} - \mathbf{v}_h\|_0 + \|\nabla(\rho - \rho_h)\|_0 \le C (h^2 \|\mathbf{v}\|_2 + h \|p\|_2)$

Implementation.

The matrix formulation reads :

$$\begin{pmatrix} \textbf{M} & \textbf{B} \\ -\textbf{B}^{\mathcal{T}} & \textbf{A} \end{pmatrix} \begin{pmatrix} \textbf{v} \\ \textbf{p} \end{pmatrix} = \begin{pmatrix} \textbf{M} \, \textbf{u} \\ \textbf{B}^{\mathcal{T}} \textbf{u} \end{pmatrix}$$

Then

$$(\mathbf{A} + \mathbf{B}^T \mathbf{M}^{-1} \mathbf{B}) \mathbf{p} = 2 \mathbf{B}^T \mathbf{u}$$

 $\mathbf{v} = \mathbf{u} - \mathbf{B} \mathbf{p}$

This is analogous to

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A stabilization of the Stokes equations

Define the spaces :

$$egin{aligned} V_h &= \{ \mathbf{v} \in \mathcal{C}^0(\overline{\Omega})^d; \ \mathbf{v}_{|T} \in \mathcal{P}_k^d, \ T \in \mathcal{T}_h, \ \mathbf{v}_{|\Gamma} = 0 \} \ Q_h &= \{ q \in \mathcal{C}^0(\overline{\Omega}); \ q_{|T} \in \mathcal{P}_\ell, \ T \in \mathcal{T}_h, \ \int_{\Omega} q = 0 \} \end{aligned}$$

We define a stabilized projection scheme by :

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$$\begin{cases} \mathbf{u}_{h}^{n+1} \in V_{h}, \ \widetilde{\mathbf{u}}_{h}^{n+1} \in V_{h}, \ p_{h}^{n+1} \in Q_{h} \\ \frac{1}{\delta t} (\widetilde{\mathbf{u}}_{h}^{n+1} - \mathbf{u}_{h}^{n}, \mathbf{v}) + (\nabla \widetilde{\mathbf{u}}_{h}^{n+1}, \nabla \mathbf{v}) = (\mathbf{f}^{n+1}, \mathbf{v}) - (\nabla p_{h}^{n}, \mathbf{v}) \qquad \mathbf{v} \in V_{h} \\ \frac{1}{\delta t} (\mathbf{u}_{h}^{n+1} - \widetilde{\mathbf{u}}_{h}^{n+1}, \mathbf{v}) + (\nabla (p_{h}^{n+1} - p_{h}^{n}), \mathbf{v}) = 0 \qquad \mathbf{v} \in V_{h} \\ - (\mathbf{u}_{h}^{n+1}, \nabla q) + \delta t (\nabla (p_{h}^{n+1} - p_{h}^{n}), \nabla q) = (\widetilde{\mathbf{u}}_{h}^{n+1}, \nabla q) \qquad q \in Q_{h} \end{cases}$$

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Convergence

(DUBOIS, TOUZANI) We take d = 2, $k = \ell = 1$. Then, under the regularity assumptions :

 $u, u_t \in L^{\infty}(H^2(\Omega)^2), \ u_{tt} \in L^{\infty}(H^1(\Omega)^2),$ $p, p_t \in L^{\infty}(H^1(\Omega)), \ p_{tt} \in L^{\infty}(L^2(\Omega)),$

we have the error bounds :

 $\begin{aligned} \|\mathbf{u} - \mathbf{u}_h\|_{\ell^{\infty}(H^1(\Omega)^2)} + \|p - p_h\|_{\ell^{\infty}(L^2(\Omega))} &\leq C(h + \delta t), \\ \|\mathbf{u} - \mathbf{u}_h\|_{\ell^{\infty}(L^2(\Omega)^2)} &\leq C(h^2 + \delta t). \end{aligned}$

This result is generalizable to the 3-D case.

Extensions

A 2nd-order scheme for Navier-Stokes equations :

- Cranck–Nicholson for the viscosity term
- Adams-Bashforth for the explicit convective term

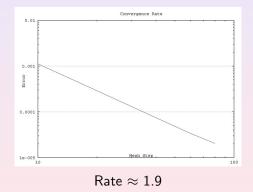
$$(\mathbf{Mv}, \mathbf{w}) := (\mathbf{v}, \mathbf{w})$$
$$(\mathbf{Kv}, \mathbf{w}) := \nu (\nabla \mathbf{v}, \nabla \mathbf{w})$$
$$(\mathbf{C}(\mathbf{v}), \mathbf{w}) := (\mathbf{v} \cdot \nabla \mathbf{v}, \mathbf{w})$$
$$(\mathbf{Bq}, \mathbf{w}) := (\nabla q, \mathbf{w})$$
$$(\mathbf{Ap}, \mathbf{q}) := (\nabla p, \nabla q)$$
$$(\mathbf{b}, \mathbf{w}) := (\mathbf{f}, \mathbf{w})$$

Mass Viscosity Convection Pressure gradient Pressure Poisson equation External forces

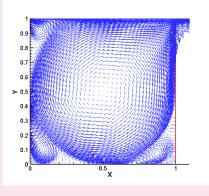
\Rightarrow 2nd-order projection scheme

$$\frac{1}{\delta t} \mathbf{M}(\widetilde{\mathbf{u}}^{n+1} - \mathbf{u}^n) + \frac{1}{2} \mathbf{K} \widetilde{\mathbf{u}}^{n+1} = \mathbf{b}^{n+\frac{1}{2}} - \mathbf{B} \mathbf{p}^n - \frac{1}{2} \mathbf{K} \mathbf{u}^n$$
$$- \frac{3}{2} \mathbf{C}(\mathbf{u}^n) + \frac{1}{2} \mathbf{C}(\mathbf{u}^{n-1})$$
$$(\mathbf{A} + \mathbf{B}\mathbf{M}^{-1}\mathbf{B}) \mathbf{q}^{n+1} = 2 \mathbf{B}^T \mathbf{u}^{n+1}$$
$$\mathbf{M} \mathbf{u}^{n+1} = \mathbf{M} \widetilde{\mathbf{u}}^{n+1} - \mathbf{B} \mathbf{q}^{n+1}$$
$$\mathbf{p}^{n+1} = \mathbf{p}^n + 2\delta t \mathbf{q}^{n+1}$$

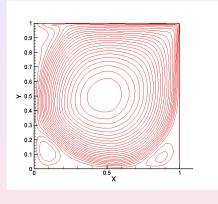
Numerical Test



Example : Driven cavity Flow



Velocity (Re=500)



Streamlines (Re=500)